

TRINITY COLLEGE FOR WOMEN NAMAKKAL Department of Physics

SOLAR ENERGY UTILIZATION 23PPHP02 - ODD Semester

Presented by

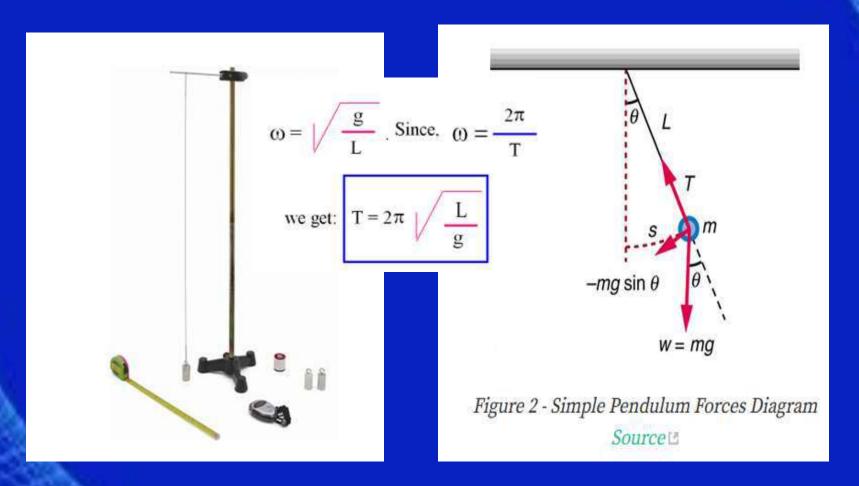
Dr.B.LAKSHMI

Head & Assistant Professor

Department of Physics

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Simple Pendulum



$$\begin{split} L &= K - V \\ L &= \frac{1}{2} m l^2 \dot{\theta}^2 - mgl(1 - \cos \theta) \\ L &= \frac{1}{2} m l^2 \dot{\theta}^2 + mgl \cos \theta - mgl \end{split}$$

Next, we plug this value into the Euler Lagrange equation, noting that $q \equiv \theta$. Thus,

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0$$

$$\frac{\partial L}{\partial \dot{\theta}} = ml^2 \dot{\theta}$$

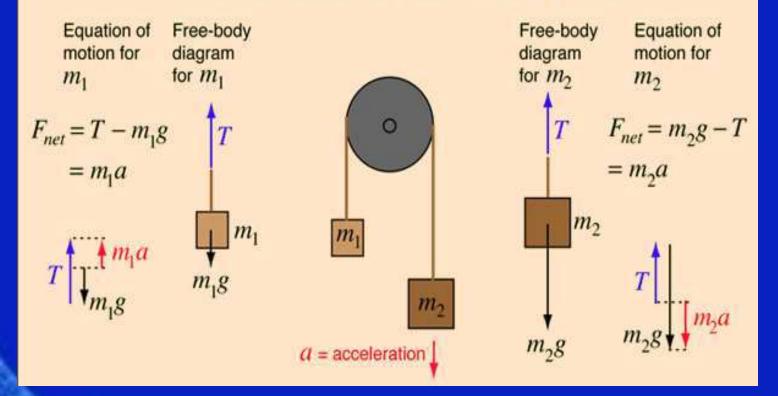
$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) = ml^2 \ddot{\theta}$$

$$\frac{\partial L}{\partial \theta} = -mgl \sin \theta$$

$$\therefore \qquad ml^2 \ddot{\theta} = -mgl \sin \theta$$
Or,
$$\ddot{\theta} = -\frac{g}{l} \sin \theta$$

Atwood's Machine

Frictionless case, neglecting pulley mass



For this idealized case the tension T is the same on both sides of the pulley. The acceleration a is the same for both masses. Solving for T gives:

$$T = m_1 g + m_1 a$$

Substituting T into the equation for m_2 gives

$$m_2 g - m_1 g - m_1 a = m_2 a$$

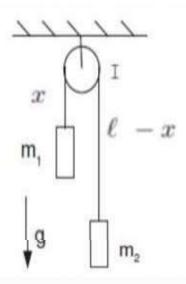
The equation of motion for the two-mass system is then:

$$(m_2 - m_1)g = (m_1 + m_2)a$$
 or $a = \frac{(m_2 - m_1)g}{(m_1 + m_2)}$

Consider the double Atwood machine. We assume that

- the pulley is light, so we can ignore their kinetic energy; and
- the rope do not slip (or they slide without friction).

The system has only one degree of freedom,



The kinetic and potential energy of the three blocks are

$$T_1 = \frac{1}{2}m_1\dot{x}^2$$

 $T_2 = \frac{1}{2}m_2 (-\dot{x})^2$
 $V_1 = -m_1gx$
 $V_2 = -m_2g(\ell - x)$

The lagrangian becomes

$$L = \frac{1}{2}(m_1 + m_2)\dot{x}^2 + m_1gx + m_2g(\ell - x)$$

Constraints:

The limitations on the motion are often called *constraints*.

➤ Holonomic Constraints:

If all constraints of the system can be expressed as equations having the form $\phi(q_1,q_2,...,q_n,t)=0$ or their equivalent, then the system is said to be *holonomic*; otherwise the system is said to be *non-holonomic*.

Example 1:

A cylinder rolling without slipping down a rough inclined plane of angle $\, \, \pmb{\alpha} \,$.

Example 2:

A horizontal cylinder of radius a rolling inside a perfectly rough hollow horizontal cylinder of radius b > a.

Example 3:

A particle constrained to move along a line under the influence of a force which is inversely proportional to the square of its distance from a fixed point and a damping force proportional to the square of the instantaneous speed.

➤ Non-Holonomic Constraints:

Any constraints that must be expressed in terms of the velocities of the particles in the system are of the form $\phi(q_1,q_2,...,q_n;\dot{q}_1,\dot{q}_2,...,\dot{q}_n;t)=0$ and constitute non-holonomic constraints unless the equations can be integrated to yield relations among the coordinates.

Example 1:

A sphere constrained to roll on a perfectly rough plane.

Example 2:

A sphere rolling down from the top of a fixed sphere.

Example 3:

A cylinder rolling [and possibly sliding] down an inclined plane of angle α .

Example 4:

A sphere rolling down another sphere which is rolling with uniform speed along a horizontal plane.

SCLERONOMIC CONSTRAINTS:

The constraints which are independent of time are called scleronomic constraints e.g. a bead sliding on a rigid curved wire fixed in space

RHEONOMIC CONSTRAINTS

The constraints which contain time explicitly are called rheonomic constraints. e.g.

- a bead sliding on a rigid curve wire moving in some prescribed fashion.
- 2)if we construct a simple pendulum whose length changes with time
- i.e. I=I(t) then the constraints expressed by the equations are time dependent, hence, rheonomic.

THANK YOU

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